

# Uniqueness and decay of correlation for Gibbs point processes, using disagreement percolation

Pierre HOUDEBERT  
Universität Potsdam

Joint work with C. Hofer-Temmel (Amsterdam)

Potsdam-Kiev Workshop, April 2019

## Definition (Configuration space)

- $\Omega$  : space of locally finite configuration  $\omega = \bigcup_{i \in I} (x_i, r_i)$ , with  $x_i \in \mathbb{R}^d$  and  $r_i \in \mathbb{R}^+$ .
- $B(\omega) = \bigcup_{(x,r) \in \omega} B(x, r)$ .
- For  $\Lambda \subseteq \mathbb{R}^d$  and  $\omega \in \Omega$ ,  $\omega_\Lambda$  is the restricted configuration of balls centered inside  $\Lambda$ ,

$$\omega_\Lambda = \omega \cap (\Lambda \times \mathbb{R}^+).$$

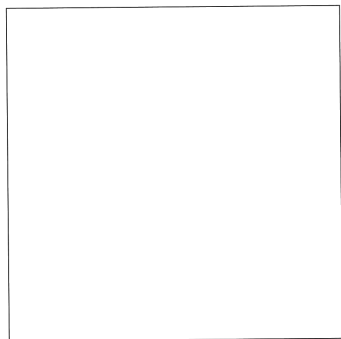
## Definition (Configuration space)

- $\Omega$  : space of locally finite configuration  $\omega = \bigcup_{i \in I} (x_i, r_i)$ , with  $x_i \in \mathbb{R}^d$  and  $r_i \in \mathbb{R}^+$ .
- $B(\omega) = \bigcup_{(x,r) \in \omega} B(x, r)$ .
- For  $\Lambda \subseteq \mathbb{R}^d$  and  $\omega \in \Omega$ ,  $\omega_\Lambda$  is the restricted configuration of balls centered inside  $\Lambda$ ,

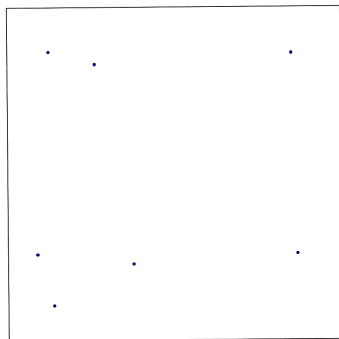
$$\omega_\Lambda = \omega \cap (\Lambda \times \mathbb{R}^+).$$

## Definition (Poisson point process)

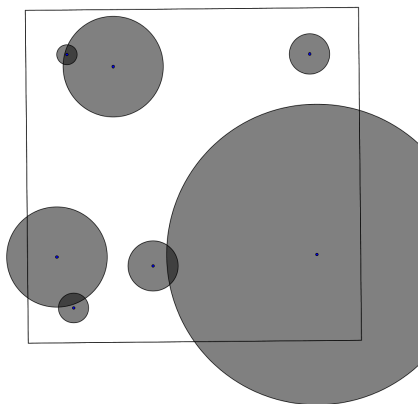
- $\pi^{z, Q}$  is the law on  $\Omega$  of a Poisson point process of intensity measure  $z \mathcal{L}^d(dx) Q(dR)$ .
- For  $\Lambda \subseteq \mathbb{R}^d$ ,  $\pi_\Lambda^{z, Q}$  is the restriction of  $\pi^{z, Q}$  on  $\Lambda \times \mathbb{R}^+$ .



- $N \sim \mathcal{P}(z|\Lambda)$ .
- $x_1, \dots, x_N$  iid uniformly on  $\Lambda$ .
- $R_1, \dots, R_N$  iid of law  $Q$ .



- $N \sim \mathcal{P}(z|\Lambda)$ .
- $x_1, \dots, x_N$  iid uniformly on  $\Lambda$ .
- $R_1, \dots, R_N$  iid of law  $Q$ .



- $N \sim \mathcal{P}(z|\Lambda)$ .
- $x_1, \dots, x_N$  iid uniformly on  $\Lambda$ .
- $R_1, \dots, R_N$  iid of law  $Q$ .

## (Infinite) Gibbs measures: definition

Let  $(H_\Lambda)$  be a fixed family of hamiltonians.

## (Infinite) Gibbs measures: definition

Let  $(H_\Lambda)$  be a fixed family of hamiltonians.

Example: "generalized" area-interaction model

$$H_\Lambda(\omega_\Lambda | \omega_{\Lambda^c}) = \mathcal{L}^d(B(\omega_\Lambda) \setminus B(\omega_{\Lambda^c})).$$



## (Infinite) Gibbs measures: definition

Let  $(H_\Lambda)$  be a fixed family of hamiltonians.

Example: "generalized" area-interaction model

$$H_\Lambda(\omega_\Lambda | \omega_{\Lambda^c}) = \mathcal{L}^d(B(\omega_\Lambda) \setminus B(\omega_{\Lambda^c})).$$

### Definition

A probability measure  $P$  belongs to the set  $\mathcal{G}(z)$  of Gibbs measures with activity  $z$  if for every bounded  $\Lambda \subseteq \mathbb{R}^d$ ,

$$P(d\omega'_\Lambda | \omega_{\Lambda^c}) = \frac{e^{-H_\Lambda(\omega'_\Lambda | \omega_{\Lambda^c})}}{\mathbf{Z}_\Lambda(\omega_{\Lambda^c})} \pi^{z, Q}(d\omega'_\Lambda), \quad (\text{DLR}_\Lambda)$$

$:= \underbrace{\mathcal{P}_{\Lambda, \omega_{\Lambda^c}}^z(d\omega'_\Lambda)}$

where  $\mathbf{Z}_\Lambda(\omega_{\Lambda^c}) = \int e^{-H_\Lambda(\omega'_\Lambda | \omega_{\Lambda^c})} \pi_{\Lambda}^{z, Q}(d\omega'_\Lambda)$ .

# (Infinite) Gibbs measures: definition

Let  $(H_\Lambda)$  be a fixed family of hamiltonians,

## Definition

A probability measure  $P$  belongs to the set  $\mathcal{G}(z)$  of Gibbs measures with activity  $z$  if for every bounded  $\Lambda \subseteq \mathbb{R}^d$ ,

$$\int f dP = \int \int f(\omega'_\Lambda \cup \omega_{\Lambda^c}) \underbrace{\frac{e^{-H_\Lambda(\omega'_\Lambda | \omega_{\Lambda^c})}}{\mathbf{Z}_{\Lambda^c}(\omega_{\Lambda^c})} \pi_\Lambda^{z, Q}(d\omega'_\Lambda)}_{:= \mathcal{P}_{\Lambda, \omega_{\Lambda^c}^c}^z(d\omega'_\Lambda)} P(d\omega), \quad (DLR_\Lambda)$$

where  $\mathbf{Z}_{\Lambda^c}(\omega_{\Lambda^c}) = \int e^{-H_\Lambda(\omega'_\Lambda | \omega_{\Lambda^c})} \pi_\Lambda^{z, Q}(d\omega'_\Lambda)$ .

## Theorem (Hofer-Temmel, H. 2017+)

*Under 3 Assumptions, we have uniqueness of the Gibbs measure for  $z$  small enough (with an explicit bound).*

*Furthermore, with one additional assumption we have exponential decay of the pair correlation function.*

- The disagreement percolation method used is coming from van den Berg & Maes (1994);
- The result was already known for the Hard-sphere model (Hofer-Temmel 2015) and for finite range Gibbs models.

## Theorem (Hofer-Temmel, H. 2017+)

*Under 3 Assumptions, we have uniqueness of the Gibbs measure for  $z$  small enough (with an explicit bound).*

*Furthermore, with one additional assumption we have exponential decay of the pair correlation function.*

- The disagreement percolation method used is coming from van den Berg & Maes (1994);
- The result was already known for the Hard-sphere model (Hofer-Temmel 2015) and for finite range Gibbs models.

### **Proof:**

Let  $P^1, P^2 \in \mathcal{G}(z)$ .

Let  $E$  be an event depending only on the configuration  $\omega_\Lambda$  for  $\Lambda$  bounded. Consider an increasing sequence  $\Lambda_n$

$$|P^1(E) - P^2(E)| \leq \int \int |\mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^1}^z(E) - \mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^2}^z(E)| P^1(d\omega^1) P^2(d\omega^2)$$

## Definition (disagreement coupling family)

A disagreement coupling family at level  $\alpha$  is a family of couplings  $\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^1, d\xi^2, d\xi^3)$  satisfying

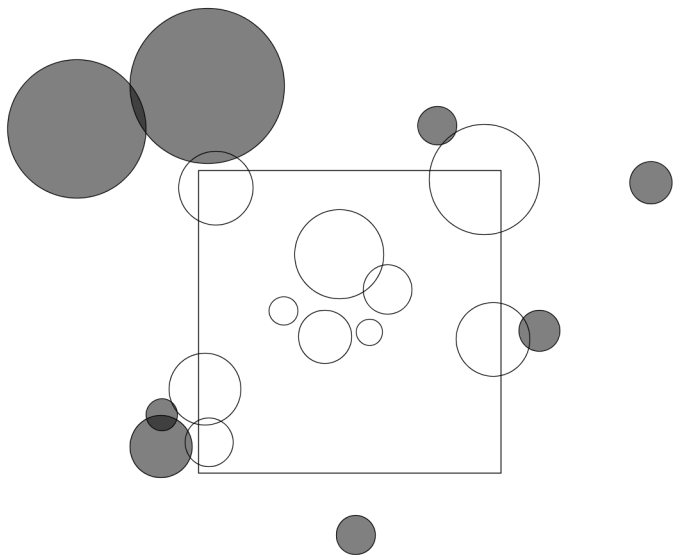
$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (\xi^1 \cup \xi^2 \subseteq \xi^3) = 1,$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^1) = \mathcal{P}_{\Lambda, \omega_{\Lambda^c}^1}^z (d\xi^1),$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^2) = \mathcal{P}_{\Lambda, \omega_{\Lambda^c}^2}^z (d\xi^2),$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^3) = \pi_{\Lambda}^{\alpha, Q} (d\xi^3),$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} \left( \forall X \in \xi^1 \Delta \xi^2 : B(X) \underset{B(\xi^3)}{\leftrightarrow} B(\omega_{\Lambda^c}^1 \cup \omega_{\Lambda^c}^2) \right) = 1.$$



Assuming the existence of a disagreement coupling family we have

$$\begin{aligned} & |\mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^1}^z(E) - \mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^2}^z(E)| \\ & \leq \int |\mathbb{1}_{\xi^1 \in E} - \mathbb{1}_{\xi^2 \in E}| \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n}^1, \omega_{\Lambda_n}^2}(d\xi^1, d\xi^2, d\xi^3) \end{aligned}$$

Theorem (Hofer-Temmel & H. 2017+)

*Under assumptions 1 and 2, there exists a disagreement coupling family at level  $\alpha$ .*

# Assumption 1 : stochastic domination

## Definition

We say that  $\tilde{P}$  stochastically dominates  $P$ , written  $P \prec \tilde{P}$ , if there exists a coupling  $\mathbf{P}$  of marginal  $P, \tilde{P}$  such that

$$\mathbf{P}(\xi^1 \subseteq \xi^2) = 1.$$



# Assumption 1 : stochastic domination

## Definition

We say that  $\tilde{P}$  stochastically dominates  $P$ , written  $P \prec \tilde{P}$ , if there exists a coupling  $\mathbf{P}$  of marginal  $P, \tilde{P}$  such that

$$\mathbf{P}(\xi^1 \subseteq \xi^2) = 1.$$

## Assumption 1

$\mathcal{P}_{\Lambda, \omega_{\Lambda^c}}^z \prec \pi_{\Lambda}^{\alpha, Q}$  for all bounded  $\Lambda$  and configurations  $\omega_{\Lambda^c}$ .  
( $\Rightarrow P \prec \pi^{\alpha, Q}$  for every  $P \in \mathcal{G}(z)$ )

# Assumption 1 : stochastic domination

## Definition

We say that  $\tilde{P}$  stochastically dominates  $P$ , written  $P \prec \tilde{P}$ , if there exists a coupling  $\mathbf{P}$  of marginal  $P, \tilde{P}$  such that

$$\mathbf{P}(\xi^1 \subseteq \xi^2) = 1.$$

## Assumption 1

$\mathcal{P}_{\Lambda, \omega_{\Lambda^c}}^z \prec \pi_{\Lambda}^{\alpha, Q}$  for all bounded  $\Lambda$  and configurations  $\omega_{\Lambda^c}$ .  
( $\Rightarrow P \prec \pi^{\alpha, Q}$  for every  $P \in \mathcal{G}(z)$ )

## Proposition (Georgii & Künetz 1997)

If  $h(X, \omega) := H_{\{X\}}(X \cup \omega) \geq C$ , then for all bounded configurations  $\omega_{\Lambda^c}$  we have

$$\mathcal{P}_{\Lambda, \omega_{\Lambda^c}}^{z, Q} \prec \pi_{\Lambda}^{ze^{-C}, Q}.$$

## Definition (disagreement coupling family)

A disagreement coupling family at level  $\alpha$  is a family of couplings  $\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^1, d\xi^2, d\xi^3)$  satisfying

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (\xi^1 \cup \xi^2 \subseteq \xi^3) = 1,$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^1) = \mathcal{P}_{\Lambda, \omega_{\Lambda^c}^1}^z (d\xi^1),$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^2) = \mathcal{P}_{\Lambda, \omega_{\Lambda^c}^2}^z (d\xi^2),$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} (d\xi^3) = \pi_{\Lambda}^{\alpha, Q} (d\xi^3),$$

$$\mathbb{P}_{\Lambda, \omega_{\Lambda^c}^1, \omega_{\Lambda^c}^2} \left( \forall X \in \xi^1 \Delta \xi^2 : B(X) \stackrel{B(\xi^3)}{\leftrightarrow} B(\omega_{\Lambda^c}^1 \cup \omega_{\Lambda^c}^2) \right) = 1.$$

## Assumption 2: "pseudo locality"

### Assumption 2

If  $B(\omega'_\Lambda) \cap B(\omega_{\Lambda^c}) = \emptyset$  then

$$H_\Lambda(\omega'_\Lambda \cup \omega_{\Lambda^c}) = H_\Lambda(\omega'_\Lambda).$$

## Assumption 2: "pseudo locality"

### Assumption 2

If  $B(\omega'_\Lambda) \cap B(\omega_{\Lambda^c}) = \emptyset$  then

$$H_\Lambda(\omega'_\Lambda \cup \omega_{\Lambda^c}) = H_\Lambda(\omega'_\Lambda).$$

### Theorem (Hofer-Temmel & H. 2017+)

*Under assumptions 1 and 2, there exists a disagreement coupling family at level  $\alpha$ .*

$$\begin{aligned}
& |\mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^1}^z(E) - \mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^2}^z(E)| \\
& \leq \int |\mathbb{1}_{\xi^1 \in E} - \mathbb{1}_{\xi^2 \in E}| \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n}^1, \omega_{\Lambda_n}^2}(d\xi^1, d\xi^2, d\xi^3)
\end{aligned}$$

$$\begin{aligned}
& |\mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^1}^z(E) - \mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^2}^z(E)| \\
& \leq \int |\mathbb{1}_{\xi^1 \in E} - \mathbb{1}_{\xi^2 \in E}| \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n}^1, \omega_{\Lambda_n}^2}(d\xi^1, d\xi^2, d\xi^3) \\
& \leq \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n}^1, \omega_{\Lambda_n}^2}(\xi_{\Lambda}^1 \Delta \xi_{\Lambda}^2 \neq \emptyset)
\end{aligned}$$

$$\begin{aligned}
& |\mathcal{P}_{\Lambda_n, \omega_{\Lambda_n^c}^1}^z(E) - \mathcal{P}_{\Lambda_n, \omega_{\Lambda_n^c}^2}^z(E)| \\
& \leq \int |\mathbb{1}_{\xi^1 \in E} - \mathbb{1}_{\xi^2 \in E}| \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n^c}^1, \omega_{\Lambda_n^c}^2}(d\xi^1, d\xi^2, d\xi^3) \\
& \leq \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n^c}^1, \omega_{\Lambda_n^c}^2}(\xi_{\Lambda}^1 \Delta \xi_{\Lambda}^2 \neq \emptyset) \\
& \leq \pi_{\Lambda_n}^{\alpha, Q} \left( \Lambda \underset{B(\xi)}{\leftrightarrow} B(\omega_{\Lambda_n^c}^1 \cup \omega_{\Lambda_n^c}^2) \right)
\end{aligned}$$



$$\begin{aligned}
& |\mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^1}^z(E) - \mathcal{P}_{\Lambda_n, \omega_{\Lambda_n}^2}^z(E)| \\
& \leq \int |\mathbb{1}_{\xi^1 \in E} - \mathbb{1}_{\xi^2 \in E}| \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n}^1, \omega_{\Lambda_n}^2}(d\xi^1, d\xi^2, d\xi^3) \\
& \leq \mathbb{P}_{\Lambda_n, \omega_{\Lambda_n}^1, \omega_{\Lambda_n}^2}(\xi_{\Lambda}^1 \Delta \xi_{\Lambda}^2 \neq \emptyset) \\
& \leq \pi_{\Lambda_n}^{\alpha, Q} \left( \Lambda \underset{B(\xi)}{\leftrightarrow} B(\omega_{\Lambda_n}^1 \cup \omega_{\Lambda_n}^2) \right) \\
& \xrightarrow{n \rightarrow \infty} 0 \text{ (with Assumption 3).}
\end{aligned}$$

## Assumption 3: existence subcritical percolation regime

### Definition

The percolation threshold of  $\pi^{z, Q}$ , written  $z_c(d, Q)$  is the critical parameter separating the subcritical percolation regime from the supercritical percolation regime.

### Assumption 1

$$\int r^d Q(dr) < \infty \quad \text{and} \quad \alpha < z_c(d, Q).$$

# Correlation function

Let  $P$  be the unique Gibbs measure (true under the 3 assumptions).

## Definition

*The pair correlation function  $\rho_2$  is defined as the density function satisfying*

$$\int \sum_{x \neq y \in \omega} F(x, y) P(d\omega) = \lambda^2 \int \int F(x, y) \rho_2(x, y) dx dy$$

# Correlation function

Let  $P$  be the unique Gibbs measure (true under the 3 assumptions).

## Definition

*The pair correlation function  $\rho_2$  is defined as the density function satisfying*

$$\int \sum_{x \neq y \in \omega} F(x, y) P(d\omega) = \lambda^2 \int \int F(x, y) \rho_2(x, y) dx dy$$

$$\rho_2(x, y) = "P(x \in \omega, y \in \omega)"$$

nb:  $x \in \omega \leftrightarrow \exists r, (x, r) \in \omega$ .

Under an additional assumption, we have exponential decay of pair correlation, i.e

$$|\rho_2(x, y) - \rho_1(x)\rho_1(y)| \leq c \times e^{-k|x-y|}.$$

Under an additional assumption, we have exponential decay of pair correlation, i.e

$$|\rho_2(x, y) - \rho_1(x)\rho_1(y)| \leq c \times e^{-k|x-y|}.$$

Assumption "exp decay"

Bounded radii:  $Q([0, r_0]) = 1$  for some fixed  $r_0$ .

# Possible generalisations

- Considering more general convex bodies instead of just balls.

# Possible generalisations

- Considering more general convex bodies instead of just balls.
- Considering a more general connection rule.



# Possible generalisations

- Considering more general convex bodies instead of just balls.
- Considering a more general connection rule.
- A disagreement coupling where the dominating measure is not a Poisson point process.  
Related to a work with D. Dereudre.
- For the exponential decay: improve the assumption?

- **Hofer-Temmel & H.:** Disagreement percolation for Gibbs ball models - 2019.
- Hofer-Temmel: Disagreement percolation for the hard-sphere model.
- van den Berg & Maes: Disagreement percolation in the study of Markov fields - 1994.
- Dereudre & H.: Sharp phase transition for the continuum Widom-Rowlinson model .
- Georgii & Küneth : Stochastic Order of Point Processes - 1997.

# Thank you for your attention